

THEORY OF GENERAL RELATIVITY

LECTURE AND TUTORIALS – PROF. DR. HAYE HINRICHSSEN / M.SC. ALEXANDRE ALVAREZ – WS 2019/20



In the tetrad basis [PxHere, public domain]

EXERCISE 8.1: LOGARITHMIC DERIVATIVE (2P)

Prove that $\Gamma^{\mu}_{\nu\mu} = \partial_{\nu} \ln \sqrt{|g|}$.

EXERCISE 8.2: TETRAD FORMALISM AND THE SPIN CONNECTION (6P)

Let $\{x^{\mu}\} = \{x^1, \dots, x^n\}$ be a coordinate system of an n -dimensional manifold \mathcal{M} and denote by $\{\partial_{\mu}\}$ and $\{dx^{\nu}\}$ the associated basis vector fields of $T\mathcal{M}$ and $T^*\mathcal{M}$, respectively. Let \mathcal{M} be equipped with the position-dependent metric $\mathbf{g} = g_{\mu\nu}(x) dx^{\mu} \otimes dx^{\nu}$.

Then, at each point of the manifold, we can introduce a set of basis vectors $\mathbf{e}_i(x)$, called *vielbein* or *tetrad*, in which the metric becomes orthonormal, i.e., it locally equals the Euclidean or Minkowski metric $\mathbf{g}(\mathbf{e}_i, \mathbf{e}_j) = \eta_{ij}$. In GR, this basis corresponds to the local free-fall basis.

At each point on \mathcal{M} we can represent the tetrad basis vector field in terms of the coordinate basis by $\mathbf{e}_i = e_i^{\mu} \partial_{\mu}$ and $\mathbf{e}^j = e^j_{\nu} dx^{\nu}$ with certain coefficients $e_i^{\mu}(x)$ and $e^j_{\nu}(x)$. Likewise there are inverse relations $\partial_{\nu} = e_{\nu}^j \mathbf{e}_j$ and $dx^{\mu} = e^{\mu}_i \mathbf{e}^i$ with coefficients $e_{\nu}^j(x)$ and $e^{\mu}_i(x)$.

- Show that the coefficients are symmetric under transposition, i.e., $e_i^{\mu}(x) = e^{\mu}_i(x)$ and $e^j_{\nu}(x) = e_{\nu}^j(x)$. (1P)
- The coefficients e^j_{ν} can be thought of as being the components of a $(1,1)$ -tensor $\mathbf{E} = e^j_{\nu} \mathbf{e}_j \otimes dx^{\nu}$. What does this form, interpreted as a linear map from the tangent space onto itself, do with an abstract tangent vector? (1P)
- Express the connection coefficients Γ^a_{ij} in the tetrad basis in terms of the connection coefficients (Christoffel symbols) $\Gamma^{\alpha}_{\mu\nu}$ in the coordinate basis. (2P)
Hint: Represent $\nabla_{\mathbf{X}}\mathbf{Y}$ in coordinate and tetrad basis and compare the components.
- True relativity geeks prefer to work with the *spin connection* ω , a Lorentz-algebra-valued 1-form defined by the components $\omega_{\mu}^{ab} := e_{\mu}^j \Gamma^a_{ij} \eta^{bi}$. Show that (2P)

$$\omega_{\mu}^{ab} = e^a_{\nu} \eta^{bi} e^{\nu}_{i;\mu} .$$

Remark: This connection can be interpreted as follows: If we go in direction \mathbf{X} , then $\omega(\mathbf{X})$ is the corresponding generator of the Lorentz transformation in the flat tetrad space that can be used to go from one tangent space to the next.

EXERCISE 8.3: GEODESIC LINE IN THE UPPER POINCARÉ HALF PLANE (4P)

Consider the upper half plane $U = \{(x, y) \mid y > 0\}$ in two dimensions with the metric

$$\mathbf{g} = y^{-2}(\mathbf{d}x \otimes \mathbf{d}x + \mathbf{d}y \otimes \mathbf{d}y) .$$

- (a) Determine the geodesic equations. (1P)
- (b) Solve the geodesic equations in arc length parameterization. You may use *Mathematica*[®] or similar tools. (2P)
- (c) What is the form of the geodesic lines? What is their maximal length? (1P)

($\Sigma = 12\mathbf{P}$)

To be handed in on Wednesday, December 11, at the beginning of the tutorial.