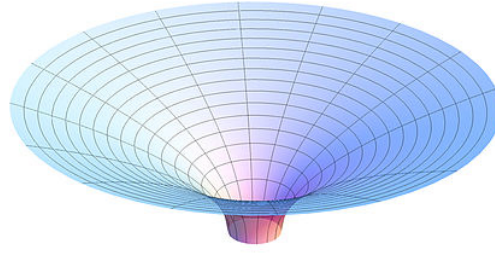


# THEORY OF GENERAL RELATIVITY

LECTURE AND TUTORIALS – PROF. DR. HAYE HINRICHSSEN / M.SC. ALEXANDRE ALVAREZ – WS 2019/20



Schwarzschild metric [Wikimedia]

## EXERCISE 10.1: PRINCIPLE OF LEAST ACTION (6P)

The electromagnetic current is a 3-form  $\mathbf{J}$  which is related to the usual charge density  $\rho$  and the current  $\vec{J} = (J_x, J_y, J_z)$  by

$$(\star\mathbf{J})^\sharp = \rho\partial_t + J_x\partial_x + J_y\partial_y + J_z\partial_z.$$

The action  $S$  of the electromagnetic field contained in a 4-volume  $G$  of the Minkowski space is given by  $S = \int_G \mathbf{L}$ , where  $\mathbf{L}$  is a 4-form called *Lagrange density*. In electrodynamics, it is given by

$$\mathbf{L} := -\frac{1}{2} d\mathbf{A} \wedge \star d\mathbf{A} + \mathbf{A} \wedge \mathbf{J}.$$

Now consider an infinitesimal variation of the electromagnetic field by  $\mathbf{A} \rightarrow \mathbf{A} + \epsilon\mathbf{Q}$ , where  $\mathbf{Q}$  is a 1-form and where  $0 < \epsilon \ll 1$  is small. The principle of least action tells us that  $S$  does not change to lowest order in  $\epsilon$ , provided that  $\mathbf{Q}$  vanishes on the boundary  $\partial G$ .

- (a) Show that the stationarity of  $S$  leads us to the condition (2P)

$$\int_G \left( -d\mathbf{Q} \wedge \star d\mathbf{A} + \mathbf{Q} \wedge \mathbf{J} \right) = 0. \quad \forall \mathbf{Q}$$

- (b) Use Stokes theorem to show that  $\int_G (d\mathbf{Q} \wedge \star d\mathbf{A}) = \int_G (\mathbf{Q} \wedge d\star d\mathbf{A})$ , using the fact that  $\mathbf{Q}|_{\partial G} = 0$ . Insert this relation into the result of (f) to show that (2P)

$$\int_G \mathbf{Q} \wedge (-d\star d\mathbf{A} + \mathbf{J}) = 0 \quad \forall \mathbf{Q} \text{ with } \mathbf{Q}|_{\partial G} = 0.$$

- (c) Show that the above result leads to the inhomogeneous Maxwell equations (1P)

$$\boxed{d\star\mathbf{F} = \mathbf{J}.}$$

- (d) Show that with the Lorenz gauge  $d^\dagger\mathbf{A} = 0$  the inhomogeneous Maxwell equations imply the wave equation  $\square\mathbf{A} = \star\mathbf{J}$ , where  $\square = dd^\dagger + d^\dagger d$  is the Laplace-de-Rham operator in Minkowski space. (1P)

**EXERCISE 10.2: TRAJECTORIES IN THE SCHWARZSCHILD METRIC (6P)**

The line element of the Schwarzschild space time in vacuum ( $T_{\mu\nu} = 0$ ) in Schwarzschild coordinates is given by

$$ds^2 = - \left(1 - \frac{2GM}{r}\right) dt^2 + \left(1 - \frac{2GM}{r}\right)^{-1} dr^2 + r^2 d\Omega^2.$$

- (a) The Lagrange function of a free particle with mass  $m$  is given by  $L = -\frac{m}{2}u^\mu u_\mu$ . Use the Euler-Lagrange equations to show that the corresponding canonical momenta  $p_t = \partial L/\partial \dot{t}$  and  $p_\phi = \partial L/\partial \dot{\phi}$  are conserved. (1P)
- (b) Show that the initial conditions  $\tau = \tau_0$ ,  $\theta(\tau_0) = \pi/2$ , and  $\dot{\theta}(\tau_0) = 0$  ensure planar motion, meaning that  $\theta = \text{const}$  for all  $\tau$ . (1P)
- (c) If the trajectory is parametrized by the proper time (eigenzeit), then  $g_{\mu\nu}\dot{x}^\mu\dot{x}^\nu = -1$ . Show that this relation can be rewritten as  $\frac{1}{2}m\dot{r}^2 + V(r) = E$  and determine  $V(r)$  and  $E$ . To what extent does  $V(r)$  differ from the corresponding potential in the Newtonian theory? Sketch both potentials qualitatively. (2P)
- (d) Consider a particle falling freely in radial direction with the initial conditions  $\tau = \tau_0$ ,  $\dot{r}(\tau_0) = \dot{\phi}(\tau_0) = \dot{\theta}(\tau_0) = \dot{t}(\tau_0) = 0$ . How much eigenzeit elapses until the particle reaches the Schwarzschild radius? (2P)

( $\Sigma = 12P$ )

To be handed in on Wednesday, January 08, at the beginning of the tutorial.